vicinity of  $\alpha \approx 10^{\circ}$ , a ratio of  $\alpha/\beta \approx 25$  is obtained, implying that the methyl twist,  $\beta$ , must be small. If the crude  $V_H(R)$  function of Bartell<sup>4</sup> is inserted into eq 3, the equation is satisfied for  $\alpha \approx 18^{\circ}$ .

If the next most important  $H \cdots H$  interactions (set B with six  $H \cdots H$  pairs) are introduced into the calculation, the angle  $\alpha$  is reduced to 16° at equilibrium, at which point the hydrogens of set A are 2.6 A apart (a distance 0.5 Å *longer* than at  $\alpha = \beta = 0^{\circ}$ ). The hydrogens of set B resist a larger twist than 16° and nestle down to a 2.6-A distance, also, at equilibrium (a distance 1.1 Å *shorter* than at  $\alpha = \beta = 0^{\circ}$ ). The remaining hydrogens play a much smaller role. In essence, then, the cooperative displacements just described have the effect of rotating the  $Si(CH<sub>3</sub>)<sub>3</sub>$  groups until the hydrogen bumps on any given group are centered on the hollows between the methyls of the adjacent groups.

According to the model, the equilibrium structure is governed more by the interhydrogen interactions than by the relatively weak bond torsional potential function,  $V(\alpha)$ . This is underscored by comparing the total effective torsional force constant for concerted rotations,  $(\partial^2 V(\alpha)/\partial \alpha^2)_{16^\circ} \approx 34 \times 10^{-12}$  ergs/radian<sup>2</sup>, including  $H \cdots H$  interactions, with the "intrinsic" contribution  $\partial^2 V_t(\alpha)/\partial \alpha^2 \approx 18 V_\alpha^0(\alpha = 16^\circ) = 1.3 \times$ 

 $10^{-12}$  ergs/radian<sup>2</sup>, excluding  $H \cdots H$  interactions. Such a large value of  $\partial^2 V(\alpha)/\partial \alpha^2$  implies that thermal amplitudes of torsional oscillations are small. **l6** The model also yields a potential energy lowering of 6.9 kcal/mol upon deformation from  $T<sub>d</sub>$  symmetry, a value much larger than thermal disordering energy. It stems almost entirely from the decrease in  $V_H(R)$  for the closest 12 hydrogens. Although the methyl-methyl interactions appreciably reduce the torsional freedom, they do not restrict the internal rotation so severely that the proton nmr signal is split.

The above experimental and calculated results are rather rough, at best. Nevertheless, they both suggest that the trimethylsilyl groups twist in a correlated manner, by perhaps 14°, in order to relieve intermethyl repulsions.

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# Electron Diffraction Determination of the Molecular Structure of Dimethylaminodichloroborane<sup>1</sup>

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Dimethylaminodichloroborane, (CH<sub>3</sub>)<sub>2</sub>NBCl<sub>2</sub>, was found to have a planar C<sub>2</sub>NBCl<sub>2</sub> skeleton. The B-N distance of 1.379  $\pm$ 0.006 **d** is the shortest B-N bond so far measured and suggests that the boron and nitrogen are linked by a double bond. The Cl-B-N and C-N-B angles were  $122.1 \pm 0.3$  and  $123.2 \pm 0.4^{\circ}$ , respectively. Root-mean-square amplitudes of vibration were determined. The structural parameters are discussed in the light of results for related molecules.

#### Introduction

Dimethylaminodichloroborane and its bromide analog were first prepared by Brown and Osthoff. $<sup>3</sup>$  The</sup> dichloride is a colorless liquid with a freezing point of  $-43^{\circ}$  and a vapor pressure of about 20 Torr at 20°. Above its melting point the compound spontaneously dimerizes to form an insoluble solid with a low vapor pressure (about 1 Torr at 91°)4 and was the first known example of self-association of an aminoborane.<sup>5</sup> The crystal structure of the cyclic dimer has been deter-

(1) This work was supported by a grant from the Horace H. Rackham School of Graduate Studies.

**(4)** H. E. Clark, Ph.D. Thesis, The University of Michigan, 1968. (5) E. L. Muetterties, Ed., "The Chemistry of Boron and Its Com-pounds," Wiley, New York, N. *Y.,* 1967, Chapter 7.

mined<sup> $\theta$ </sup> and the monomer has been studied in the infrared and Raman regions by Whitmer.7

An electron diffraction study of monomeric  $(CH_3)_2$ - $NBCl<sub>2</sub>$  was undertaken to investigate the character of the B-N bond.

#### Experimental Section

Dimethylaminodichloroborane was prepared by the procedure of Brown and Osthoff<sup>3</sup> and Whitmer,<sup>7</sup> and the purity was checked by infrared spectroscopy. Except while being used, the sample container was stored at  $-78^{\circ}$  to prevent formation of the unwanted dimer.

The diffraction apparatus and the procedure for obtaining the experimental leveled intensity from the 21- and 11-cm camera

<sup>(16)</sup> Hence it seems reasonable to interpret the results of Figure 3 in terms of a torsional deformation instead of "torsional shrinkage effects" of the sort discussed by L. S. Bartell and D. A. Kohl, *J. Chem.* Phys., **39,** 3097 (1963).

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<sup>(3)</sup> C. A. Brown and R. C. Osthoff, *J. Amev. Chem. SOC.,* **74, 2340** (1952).

<sup>(6)</sup> H. Hess, *2. Rvistallogv., Kvistallgeometvie, Kvistollphrs., Kvislall chem.,* **118,** 361 (1963).

**<sup>(7)</sup>** J. C. Whitmer, Ph.D. Thesis, The University of Michigan, 1965.



Figure 1. $-M$ olecular intensity curves for  $(CH_3)_2NBC1_2$ :  $\Delta sM(s)$  $= sM(s)_{\text{exptl}} - sM(s)_{\text{calod}}.$ 

range data (taken with an  $r^3$  sector) have been described else $where.$ <sup>8, 9</sup>

## Analysis **of** Data

Leveled experimental intensities were converted to reduced molecular intensities and fitted by a theoretical reduced intensity function according to a least-squares procedure. Experimental intensity curves were interpolated to convenient values of the scattering variable, and the resultant smoothing was taken into account in error analysis. Comparisons between experimental and theoretical points were carried out using a weighting function

$$
W(s) = C - \exp\{-\alpha [(10s/\pi) - 10]^2\} \tag{1}
$$

with  $\alpha = 0.005$  and  $C = 1.2$  for merged<sup>10</sup> 21- and 11cm camera range data.

Experimental radial distribution functions were calculated by the usual procedure $^{9,11-14}$  with theoretical data from  $s = 0$  to  $4.5 \text{ Å}^{-1}$  blended into experimental data from  $s = 3.5$  to 37 Å<sup>-1</sup>. The value of *b* in the damping factor  $e^{-bs^2}$  was taken to be 0.00125.

Asymmetry constants  $a$  were estimated<sup>15</sup> to be 2.5, 1.9, 1.8, and 1.6  $\AA^{-1}$  for the H-C, B-N, C-N, and B-C1 bonded distances, respectively, and were taken to be

- (9) L. S. Bartell, D. **A.** Kohl, B. L. Carroll, and R. M. Gavin, Jr., *ibid.,*  **43,** 3079 (1985); R. **A.** Bonham and L. S. Bartell, *ibid.,* **31,** 702 (1959).
- **(10)** Blended in the overlap region as in constructing the radial distribution function.
- (11) K. Kuchitsu and L. *S.* Bartell, *J. Chem. Phys.,* **36,** 1945 (1961).
- (12) R. **A.** Bonhamand L. S. Bartell, *J. Amev.* Chem. *SOL.,* **81,** 3491 (1959). (13) J. **A.** Ibers and J. **A.** Hoerni, *Acta Cvystallogv., 7,* 405 (1954).
- (14) **I,.** *S.* Bartell and T. L. Boates, to be submitted for publication.
- (15) D. R. Herschbach and V. **W.** Laurie, *J.* Chem. *Phys.,* **35,** 458 (1961).



Figure 2.—Radial distribution function for  $(CH_3)_2NBC1_2$ :  $\Delta f(r)$ <br>=  $f(r)_{\text{exptl}} - f(r)_{\text{oated}}$ .

TABLE I



<sup>a</sup> Distances in ångströms. Angles in degrees. Estimated uncertainties include the effects of known systematic errors and random errors inferred from least-squares analyses according to L. S. Bartell in "Physical Methods in Chemistry," **A.** Weissberger and B. W. Rossiter, Ed., 4th ed, Interscience, New York, N. Y., in press. The interaction between random and systematic errors as expressed in terms of **eq** 27 in this reference has not, however, been included. Hence, the standard errors listed are probably overoptimistic, particularly for the angles. *See,* for example, Figure 3. <sup>b</sup> Angle between the planes C-N-B and Cl<sub>2</sub>-BN.  $\mathfrak{e} \, l_g(B \cdots C)$  constrained to equal  $l_g(C \cdots C) = \mathfrak{e} \, l_g$ - $(C \cdots C)_{short}$  and  $l_g(C \cdots C)_{long}$  constrained to equal  $l_g(N \cdots C)$ plus 0.010 and 0.012 **A,** respectively.

 $1.0 \text{ Å}^{-1}$  for all nonbonded distances. No correction was made for the Bastiansen-Morino shrinkage effect.<sup>16</sup>

### Results

Figure 1 shows the molecular intensity curve<sup>10</sup> determined for  $(CH_3)_2NBC1_2$ . The mean index of resolution  $R = M(s)_{\text{expt}}/M(s)_{\text{caled}}$  was 1.04 for both the 21and 11-cm camera ranges, The experimental radial distribution function is illustrated in Figure *2.* The (16) *Y.* Morino, S. J. Cyvin, K. Kuchitsu, and T. Iijima, *ibid.,* **36,** <sup>1109</sup> (1962).

<sup>(8)</sup> L. *S.* Bartell, K. Kuchitsu, and R. J. DeNeui, *J.* Chem. *Phys.,* 86, 1211 (1961).



TABLE I1

<sup>a</sup> Values are  $\times 10$ 4. Based on 109 intensity values interpolated from 278 data points. Units for distances and amplitudes are ångstroms, and for angles they are radians; the index of resolution *R* is dimensionless. Matrix elements are given by  $\sigma_{ij} = \text{sign}[B_{ij}^{-1}]$ .  $\{B_{ij}-1|\mathbf{v}'\mathbf{w}\mathbf{v}/(n-m)\}^{1/2}$ , where the notation corresponds to that of O. Bastiansen, L. Hedberg, and K. Hedberg, *J. Chem. Phys.*, 27, 1311 (1957). See also K. Hedberg and M. Iwasaki, *Acta Crystallogr.,* **17,** 529 (1964). Since the elements are based on a diagonal (nonoptimum) weight matrix, they do not represent *bona fide* standard errors. \* Angle between the planes C-N-B and the plane C12BN.  $\partial^p$  *I(*  $B \cdots C$ ) was constrained to equal  $l(C \cdots C)$ .  $d$   $l(C \cdots C)$ <sub>bhort</sub> and  $l(C \cdots C)$ <sub>long</sub> were constrained to equal  $l(N \cdots C)$  plus 0.010 and 0.012 **A,** respectively.

TABLE I11



<sup>a</sup> The 21-cm camera data are given above, and the 11-cm data, below. The function  $M(s)$  is given by  $[(I_0/B) - 1]$ .

results of our investigation are summarized in Table I. Experimental data were analyzed by least-squares fittings of the intensity for each camera distance and for of the error matrix<sup>17</sup> computed during the final leasta blend of the two camera distances. **A** representation possibility of any appreciable twist about the BN bond squares runs on the blended intensity may be seen in Table 11, **A** listing of the experimental leveled intensity and the background used in data analysis as funccamera distances is given in Table 111. tions of the scattering variable for the 21- and 11-cm to what extent models with slightly pyramidal nitrogen

**(1965);** Y. **Murata and Y.** Morino, *ibid.,* **20, 605 (1966). (17) Y.** Morino, **K. Kuchitsu, and Y. Murata, Acta** *C~ystallogr.,* **18, 549** 

## **Discussion**

The experimental radial distribution excludes the and is completely consistent with a planar  $C_{2v}$  structure analogous to that of a substituted ethylene. Although it was unnecessary to use a nonplanar skeleton model to fit the diffraction data, an effort was made to establish atoms are compatible with the data. The necessary calculations were performed for a series of  $C<sub>s</sub>$  models with nonplanar skeletons in which both methyl groups



TABLE IV

<sup>a</sup> Distances in ångströms; angles in degrees. <sup>b</sup> This work. <sup>c</sup> See ref 18. <sup>d</sup> See ref 20. <sup>e</sup> S. Konaka, Y. Murata, K. Kuchitsu, and Y. Morino, Bull. Chem. Soc. Jap., 39, 1134 (1966). / K. P. Coffin and S. H. Bauer, J. Phys. Chem., 59, 193 (1955).

were displaced from the Cl<sub>2</sub>BN plane. It is evident from Figure 3 and the error matrix that the answer based on electron diffraction data is indeterminate; distortions less than about 10° cannot be ruled out.



Figure 3.—Minima in standard deviation curve vs. angle between the planes C-N-B and the plane  $Cl<sub>2</sub>BN$  (C<sub>s</sub> symmetry  $model)$  for  $(CH<sub>3</sub>)<sub>2</sub>NBCl<sub>2</sub>$ .

Several compounds related to  $(CH_3)_2NBCl_2$  are compared in Table IV. The most striking feature is the very large  $(0.1 \pm 0.02 \text{ Å})$  difference between the B-N bond lengths in  $(CH_3)_2NBCl_2$  and  $(SiH_3)_2NBF_2$ .<sup>18</sup> One way of rationalizing this difference is in terms of the postulate of Robiette, et al., that substantial  $(p \rightarrow d)\pi$ bonding from nitrogen to silicon occurs in the latter compound, consistent with the known planarity of  $N(SiH<sub>3</sub>)<sub>3</sub>$ <sup>19</sup> An analogous delocalization of the nitrogen lone pair along the C-N bonds in  $(CH_3)_2NBC1_2$ , of course, is not favorable. The existence of  $(p \rightarrow d)\pi$ bonding in the silicon compound would presumably tend to shorten the Si-N bond and lengthen the B-N bond. From the data collected in Tables IV and V it appears that the B-N bond in  $(CH_3)_2NBCl_2$  is within the range expected for a normal B-N double bond, and

the 1.48  $\pm$  0.02 Å B-N bond reported<sup>18</sup> for  $(SiH<sub>3</sub>)<sub>2</sub>$  $NBF<sub>2</sub>$  is intermediate between a single and a double bond.





<sup>a</sup> L. S. Bartell and H. K. Higginbotham, J. Chem. Phys., 42, 851 (1965). <sup>b</sup> L. S. Bartell, E. A. Roth, C. D. Hollowell, K. Kuchitsu, and J. E. Young, Jr., ibid., 42, 2683 (1965). . I. L. Karle, ibid., 20, 65 (1952). <sup>d</sup> P. H. Clippard, Ph.D. Thesis, University of Michigan, 1969. B-N bond determined for adducts  $(CH_3)_3NBCl_3$ ,  $(CH_3)_3NBBr_3$ , and  $(CH_3)_3NBI_3$ . See also P. H. Clippard, R. C. Taylor, and J. C. Hanson, Inorg. Chem., in press; L. E. Sutton, "Tables of Internuclear Distances and Configurations in Molecules and Ions," Special Publications No. 11 and 18, The Chemical Society, London, 1958 and 1965. The higher (Clippard) values are the most accurate but may be unrepresentatively long because of steric deformation. *•* This work. *' S. H. Bauer, J. Amer. Chem. Soc.*, 60, 524 (1938); W. Harshbarger, G. Lee, R. F. Porter, and S. H. Bauer, Inorg. Chem., 8, 1683 (1969).

The molecule  $(CH_3)_2NNO_2$  recently studied by Stölevik and Rademacher<sup>20</sup> resembles  $(CH_3)_2NBC1_2$ , being isoelectronic with  $(CH_3)_2NBF_2$ . No tendency of the bonds around the trivalent nitrogen to deviate from a planar configuration was detected.

In summary, the B-N bond length determined in this work,  $1.379 \pm 0.006$  Å, is the shortest boron-nitrogen distance yet measured. The bond length, the apparent skeletal planarity, and the spontaneous dimerization exhibited by the molecule are all consistent with the characterization of the B-N bond as a double bond. Values of comparable bonds in the isoelectronic C-C and B-N series are listed in Table V, where a marked similarity between the series is apparent.

<sup>(18)</sup> A. G. Robiette, G. M. Sheldrick, W. S. Sheldrick, B. Beagley, D. W. J. Cruickshank, J. J. Monaghan, B. J. Aylett, and I. A. Ellis, Chem. Commun., 909 (1968).

<sup>(19)</sup> K. Hedberg, J. Amer. Chem. Soc., 77, 6491 (1955).

<sup>(20)</sup> P. Rademacher and R. Stölevik, Acta Chem. Scand., 23, 660 (1969); R. Stölevik and P. Rademacher, ibid., 23, 672 (1969).